Toward less synchronous composable multilevel methods for implicit multiphysics simulation

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Outline

Multiphysics and methods

Coupling software in PETSc

Hardware and consequences

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Multiphysics problems

Examples

- Saddle-point problems (e.g. incompressibility, contact)
- Stiff waves (e.g. low-Mach combustion)
- Mixed type (e.g. radiation hydrodynamics, ALE free-surface flows)
- Multi-domain problems (e.g. fluid-structure interaction)
- Full space PDE-constrained optimization

Software/algorithmic considerations

- Separate groups develop different "physics" components
- Do not know a priori which methods will have good algorithmic properties
- Achieving high throughput is more complicated
- Multiple time and/or spatial scales
 - Splitting methods are delicate, often not in asymptotic regime
 - ► Strongest nonlinearities usually non-stiff: prefer explicit for TVD limiters/shocks

The Great Solver Schism: Monolithic or Split?

Monolithic

- Direct solvers
- Coupled Schwarz
- Coupled Neumann-Neumann (need unassembled matrices)
- Coupled multigrid
- X Need to understand local spectral and compatibility properties of the coupled system

Split

- Physics-split Schwarz (based on relaxation)
- Physics-split Schur (based on factorization)
 - approximate commutators SIMPLE, PCD, LSC
 - segregated smoothers
 - Augmented Lagrangian
 - "parabolization" for stiff waves

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- X Need to understand global coupling strengths
- Preferred data structures depend on which method is used.
- Interplay with geometric multigrid.

Splitting for Multiphysics

$$\begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} f \\ g \end{bmatrix}$$

► Relaxation: -pc_fieldsplit_type [additive,multiplicative,symmetric_multiplicative] $\begin{bmatrix} A \\ D \end{bmatrix}^{-1} \begin{bmatrix} A \\ C \end{bmatrix}^{-1} \begin{bmatrix} A \\ 1 \end{bmatrix}^{-1} \begin{pmatrix} A \\ 1 \end{bmatrix}^{-1} \begin{pmatrix} A \\ D \end{bmatrix}^{-1} \begin{bmatrix} A \\ C \end{bmatrix}^{-1} \begin{pmatrix} A \\ D \end{pmatrix}^{-1} \begin{pmatrix} A \\$

Gauss-Seidel inspired, works when fields are loosely coupled
 Factorization: -pc_fieldsplit_type schur

$$\begin{bmatrix} A & B \\ & S \end{bmatrix}^{-1} \begin{bmatrix} 1 \\ CA^{-1} & 1 \end{bmatrix}^{-1}, \qquad S = D - CA^{-1}B$$

- robust (exact factorization), can often drop lower block
- how to precondition S which is usually dense?
 - interpret as differential operators, use approximate commutators

How much nesting?

$$P_{1} = \begin{pmatrix} J_{uu} & J_{up} & J_{uE} \\ 0 & B_{pp} & 0 \\ 0 & 0 & J_{EE} \end{pmatrix}$$

- *B_{pp}* is a mass matrix in the pressure space weighted by inverse of kinematic viscosity.
- Elman, Mihajlović, Wathen, JCP 2011 for non-dimensional isoviscous Boussinesq.
- Works well for non-dimensional problems on the cube, not for realistic parameters.

$$P = \begin{bmatrix} \begin{pmatrix} J_{uu} & J_{up} \\ J_{pu} & 0 \end{pmatrix} & \\ \begin{pmatrix} J_{Eu} & J_{Ep} \end{pmatrix} & J_{EE} \end{bmatrix}$$

- Inexact inner solve using upper-triangular with B_{pp} for Schur.
- Another level of nesting.
- GCR tolerant of inexact inner solves.
- Outer converges in 1 or 2 iterations.
- Low-order preconditioning full-accuracy unassembled high order operator.
- Build these on command line with PETSc PCFieldSplit.

Example 3×3 problem with nested 2×2 split

```
-fieldsplit_s_ksp_type gcr
-fieldsplit_s_ksp_rtol 1e-1
-fieldsplit_s_ksp_monitor_vht
-fieldsplit_s_ksp_monitor_singular_value
-fieldsplit_s_pc_type fieldsplit
-fieldsplit_s_pc_fieldsplit_type schur
-fieldsplit_s_pc_fieldsplit_real_diagonal
-fieldsplit_s_pc_fieldsplit_schur_factorization_type lower
-fieldsplit_s_fieldsplit_u_ksp_type gmres
-fieldsplit_s_fieldsplit_u_ksp_max_it 10
-fieldsplit_s_fieldsplit_u_pc_type asm
-fieldsplit_s_fieldsplit_u_sub_pc_type ilu
-fieldsplit_s_fieldsplit_u_sub_pc_factor_levels 1
-fieldsplit_s_fieldsplit_u_ksp_converged_reason
-fieldsplit_s_fieldsplit_p_ksp_type preonly
-fieldsplit_s_fieldsplit_p_ksp_max_it 1
-fieldsplit_s_fieldsplit_p_pc_type jacobi
-fieldsplit_e_ksp_type gmres
-fieldsplit_e_ksp_converged_reason
-fieldsplit_e_pc_type asm
-fieldsplit_e_sub_pc_type ilu
-fieldsplit_e_sub_pc_factor_levels 2
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Monolithic approaches

Parallel direct solver

-dm_mat_type aij -pc_type lu -pc_factor_mat_solver_package mumps

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Coupled nonlinear multigrid accelerated by NGMRES with multi-stage smoothers

-lidvelocity 200 -grashof 1e4 -snes_grid_sequence 5 -snes_monitor -snes_view -snes_type ngmres -npc_snes_type fas -npc_snes_max_it 1 -npc_fas_coarse_snes_type ls -npc_fas_coarse_ksp_type preonly -npc_fas_snes_type ms -npc_fas_snes_max_it 1 -npc_fas_ksp_type preonly -npc_fas_pc_type pbjacobi -npc_fas_snes_ms_type m62 -npc_fas_snes_max_it 1

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- package each "physics" independently
- solve single-physics and coupled problems
- semi-implicit and fully implicit
- reuse residual and Jacobian evaluation unmodified
- direct solvers, fieldsplit inside multigrid, multigrid inside fieldsplit without recompilation
- use the best possible matrix format for each physics (e.g. symmetric block size 3)

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- matrix-free anywhere
- multiple levels of nesting

MomentumStokes Pressure

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Boundary Layer

Ocean

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Work in Split Local space, matrix data structures reside in any space.

MatGetLocalSubMatrix(Mat A,IS rows,IS cols,Mat *B);

- Primarily for assembly
 - B is not guaranteed to implement MatMult
 - The communicator for B is not specified, only safe to use non-collective ops (unless you check)
- IS represents an index set, includes a block size and communicator
- MatSetValuesBlockedLocal() is implemented
- MatNest returns nested submatrix, no-copy
- No-copy for Neumann-Neumann formats (unassembled across procs, e.g. BDDC, FETI-DP)
- Most other matrices return a lightweight proxy Mat
 - COMM_SELF
 - Values not copied, does not implement MatMult
 - Translates indices to the language of the parent matrix

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Multiple levels of nesting are flattened

Stokes + Implicit Free Surface

$$\begin{bmatrix} \eta D_{ij}(\boldsymbol{u}) \end{bmatrix}_{,j} - p_{,i} = f_i$$
$$u_{k,k} = 0$$
$$\hat{x}_i = \hat{x}_i^{t-\Delta t} + \Delta t \, u_i(\hat{x}_i)$$



COORDINATE RESIDUALS

$$F_x := -u_i + \frac{\hat{x}_i}{\Delta t} - \frac{\hat{x}_i^{t-\Delta}}{\Delta t}$$

[We use a full Lagrangian update of our mesh, with no remeshing]



May, Le Pourhiet & Brown: Coupled Geodynamics



"Drunken seaman", Rayleigh Taylor instability test case from Kaus et al., 2010. Dense, viscous material (yellow) overlying less dense, less viscous material (blue).



16



Stokes + Implicit Free Surface



* The nonlinear residual ALWAYS increases from one step to the next.

* A nonlinear solve is required to control the error.

* An accurate nonlinear solve on the first time step, combined with 1 or 2 nonlinear iterations on subsequent steps still results in severe errors. *This is true even when dt is small.*

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18

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On-node hardware roadmap

Hardware trends

- More cores (keep hearing $\mathcal{O}(1000)$ per node)
- Long vector registers (32B for AVX and BG/Q, 64B for MIC)
- Must use SMT to hide memory latency (POWER7)
- Must use SMT for floating point performance (GPU, BG/Q)
- Large penalty for non-contiguous memory access

"Free flops", but how can we use them?

- High order methods good: better accuracy per storage
- High order methods bad: work unit gets larger
- GPU threads have very little memory, must keep work unit small
- Need library composability, keep user contribution embarrassingly parallel

How to program this beast?

- Decouple physics from discretization
 - Expose small, embarrassingly parallel operations to user
 - Library schedules user threads for reuse between kernels
 - User provides physics in kernels run at each quadrature point
 - Continuous weak form: find $u \in \mathscr{V}_D$

$$v^T F(u) \sim \int_{\Omega} v \cdot f_0(u, \nabla u) + \nabla v : f_1(u, \nabla u) = 0, \quad \forall v \in \mathscr{V}_0$$

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- Similar form at faces, may involve Riemann solve
- Library manages reductions
 - Interpolation and differentiation on elements
 - Exploit tensor product structure to keep working set small
 - Assembly into solution/residual vector (sum over elements)

Nodal hp-version finite element methods



1D reference element

- Lagrange interpolants on Legendre-Gauss-Lobatto points
- Quadrature \hat{R} , weights \hat{W}
- Evaluation: \hat{B}, \hat{D}

3D reference element

$$\begin{array}{ll} \hat{W} = \hat{W} \otimes \hat{W} \otimes \hat{W} & \hat{D}_0 = \hat{D} \otimes \hat{B} \otimes \hat{B} \\ \hat{B} = \hat{B} \otimes \hat{B} \otimes \hat{B} & \hat{D}_1 = \hat{B} \otimes \hat{D} \otimes \hat{B} \\ \hat{D}_2 = \hat{B} \otimes \hat{B} \otimes \hat{D} \end{array}$$

These tensor product operations are very efficient, 70% of peak flop/s

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Operations on physical elements

Mapping to physical space

$$x^e: \hat{K} \to K^e, \quad J^e_{ij} = \partial x^e_i / \partial \hat{x}_j, \quad (J^e)^{-1} = \partial \hat{x} / \partial x^e$$

Element operations in physical space

$$B^{e} = \hat{B} \qquad W^{e} = \hat{W}\Lambda(|J^{e}(r)|)$$
$$D_{i}^{e} = \Lambda\left(\frac{\partial \hat{x}_{0}}{\partial x_{i}}\right)\hat{D}_{0} + \Lambda\left(\frac{\partial \hat{x}_{1}}{\partial x_{i}}\right)\hat{D}_{1} + \Lambda\left(\frac{\partial \hat{x}_{2}}{\partial x_{i}}\right)\hat{D}_{2}$$
$$(D_{i}^{e})^{T} = \hat{D}_{0}^{T}\Lambda\left(\frac{\partial \hat{x}_{0}}{\partial x_{i}}\right) + \hat{D}_{1}^{T}\Lambda\left(\frac{\partial \hat{x}_{1}}{\partial x_{i}}\right) + \hat{D}_{2}^{T}\Lambda\left(\frac{\partial \hat{x}_{2}}{\partial x_{i}}\right)$$

Global problem is defined by assembly

$$F(u) = \sum_{e} \mathscr{E}_{e}^{T} \left[(B^{e})^{T} W^{e} \Lambda(f_{0}(u^{e}, \nabla u^{e})) + \sum_{i=0}^{d} (D_{i}^{e})^{T} W^{e} \Lambda(f_{1,i}(u^{e}, \nabla u^{e})) \right] = 0$$

where $u^e = B^e \mathscr{E}^e u$ and $\nabla u^e = \{D^e_i \mathscr{E}^e u\}_{i=0}^2$

Representation of Jacobians, Automation

- For unassembled representations, decomposition, and assembly
- Continuous weak form: find u

$$v^T F(u) \sim \int_{\Omega} v \cdot f_0(u, \nabla u) + \nabla v : f_1(u, \nabla u) = 0, \quad \forall v \in \mathscr{V}_0$$

Weak form of the Jacobian J(u): find w

- Terms in $[f_{i,j}]$ easy to compute symbolically, AD more scalable.
- ► Nonlinear terms *f*₀,*f*₁ usually have the most expensive nonlinearities in the computation of scalar material parameters
 - Equations of state, effective viscosity, "star" region in Riemann solve
 - Compute gradient with reverse-mode, store at quadrature points.
 - Perturb scalars, then use forward-mode to complete the Jacobian.

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Flip for action of the adjoint.

Conservative (non-Boussinesq) two-phase ice flow

Find momentum density ρu , pressure p, and total energy density E:

$$(\rho u)_t + \nabla \cdot (\rho u \otimes u - \eta D u_i + p 1) - \rho g = 0$$

$$\rho_t + \nabla \cdot \rho u = 0$$

$$E_t + \nabla \cdot ((E+p)u - k_T \nabla T - k_\omega \nabla \omega) - \eta D u_i : D u_i - \rho u \cdot g = 0$$

- Solve for density ρ, ice velocity u_i, temperature T, and melt fraction ω using constitutive relations.
 - Simplified constitutive relations can be solved explicitly.
 - Temperature, moisture, and strain-rate dependent rheology η.

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- High order FEM, typically Q₃ momentum & energy
- DAEs solved implicitly after semidiscretizing in space.
- Preconditioning using nested fieldsplit

Performance of assembled versus unassembled



- High order Jacobian stored unassembled using coefficients at quadrature points, can use local AD
- Choose approximation order at run-time, independent for each field
- Precondition high order using assembled lowest order method
- Implementation > 70% of FPU peak, SpMV bandwidth wall < 4%</p>

Memory Bandwidth

Operation		Arithmetic Intensity (flops per byte)	
Sparse matrix-vector product		1/6	
Dense matrix-vector product		1/4	
Unassembled matrix-vector product		pprox 8	
High-order residual evaluation		> 5	
Processor	BW (GB/s)	Peak (GF/s)	Balanced AI (F/B)
Sandy Bridge 6-core	21*	150	7.2
Magny Cours 16-core	42*	281	6.7
Blue Gene/Q node	43	205	4.8
GeForce 9400M	21	54	2.6
GTX 285	159	1062	6.8
Tesla M2050	144	1030	7.1

Prospects for reducing synchronization

- Dot products and norms
 - orthogonality is a powerful concept
 - dot product/norm fusion in CG variants
 - Iatency-tolerant Krylov methods, TSQR for GMRES
 - nonlinear methods (e.g. NGMRES, BFGS, line searches)
 - hierarchical methods to limit system-wide norms
 - setting up smoothers and coarsening rates for AMG
- additive coarse grids
- subphysics on subcommunicators, even within multigrid context
- s-step methods (and other fusion)
 - often spoiled by algorithmic requirements of preconditioning

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- relevant for multigrid smoothers
- difficult crossovers for 3D problems

Multigrid is always strong scaling

- Finest level is chosen by the application (might have big subdomains)
- All coarsened levels choose communicator size based on strong scaling limit
- Optimizing the strong scaling limit pays off consistently
- Rapid coarsening is important (2:1 semi-coarsening not okay any more)

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Advanced Analysis

- Uncertainty quantification
 - intrusive vs unintrusive methods, multilevel
 - uncertainty in modeling error
 - use subdifferentials for non-smooth processes
 - unified handling of heterogeneous observational data
- PDE-constrained optimization
 - multi-objective
 - robustness
 - rich problem description
 - fusing algorithmic steps (LNK and coupled DD fuse gradients with progress)

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- Exploring stability manifolds
 - solving bordered linear and nonlinear systems
- nearly time-periodic nonlinear problems
 - identifying cycles in ocean models, turbomachinery

Software challenges

- Which interfaces do users have to interact with?
 - "F"ramework vs library
 - Extensibility is critical for multiphysics
- Asynchronous interfaces crossing module boundaries
 - How to ensure progress?
- Merge communication on multiple levels or between multiple physics
- Fusing coarse level operations
- Working with non-nested communicators is tricky
- Current solutions for hierarchical memory are bad for libraries
 - I want a communicator-like object
 - I want a way to allocate memory explicitly/relative to algorithmic dependencies instead of implicit "first touch"
- Time integration: IMEX, multirate, parallel in time
 - method of lines: $g(\dot{u}, u, t) = f(u, t)$
 - Lax-Wendroff time integration is harder for composable software